

# Theoretical study on the encapsulation of Pd<sub>3</sub>-based transition metal clusters inside boron nitride nanotubes

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Received: 30 August 2012 / Accepted: 22 October 2012 / Published online: 13 November 2012  
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**Abstract** Chemical functionalization of the boron nitride nanotube (BNNT) allows a wider flexibility in engineering its electronic and magnetic properties as well as chemical reactivity, thus making it have potential applications in many fields. In the present work, the encapsulation of 13 different Pd<sub>3</sub>M (*M*=Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn, Pd, Pt, and Au) clusters inside the (10, 0) BNNT has been studied by performing comprehensive density functional theory (DFT) calculations. Particular attention is paid to searching for the stable configurations, calculating the corresponding binding energies, and evaluating the effects of the encapsulation of Pd<sub>3</sub>M cluster on the electronic and magnetic properties of BNNT. The results indicate that all the studied Pd<sub>3</sub>M clusters can be stably encapsulated inside the (10, 0) BNNT, with binding energies ranging from −0.96 (for Pd<sub>3</sub>Sc) to −5.31 eV (for Pd<sub>3</sub>V). Moreover, due to a certain amount of charge transfer from Pd<sub>3</sub>M clusters to BNNT, certain impurity states are induced within the band gap of pristine BNNT, leading to the reduction of the band gap in various ways. Most Pd<sub>3</sub>M@BNNT nanocomposites exhibit nonzero magnetic moments, which mainly originate from the contribution of the Pd<sub>3</sub>M clusters. In particular, the adsorption of O<sub>2</sub> molecule on BNNT is greatly enhanced due to Pd<sub>3</sub>M encapsulation. The elongation of O–O bonds of the adsorbed O<sub>2</sub> molecules indicates that Pd<sub>3</sub>M@BNNT could be used to fabricate the oxidative catalysis.

**Keywords** Boron nitride nanotube · DFT · Encapsulation · Pd<sub>3</sub>M cluster

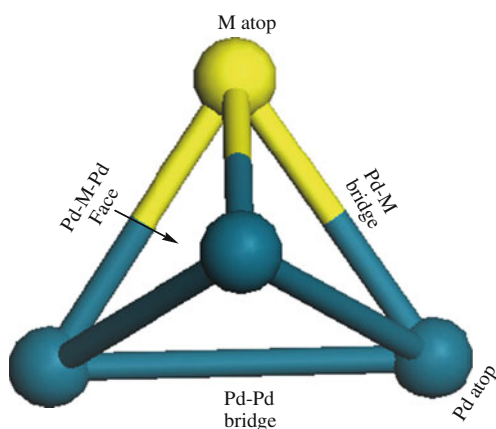
## Introduction

Since its discovery in 1991 [1], the carbon nanotube (CNT) has attracted increasing attention due to its potential applications. The impressive progress in CNT research has motivated scientists to explore other 1D atomic based materials. Among them, the boron nitride nanotube (BNNT) [2] has become a hotly pursued system, as it shares the same honeycomb lattice structure as CNT. Although BNNT has similar geometry as CNT, it may have more important advantages over CNT. For example, BNNTs not only have high thermal conductivity but also high oxidation resistivity as well as high thermal and chemical stability [3, 4]. These properties render them good candidates for developing nanotube-based devices in hazardous and high-temperature environments.

Moreover, BNNTs are semiconductors with a wide band gap (~5.5 eV) that are weakly dependent on the tube diameter, chirality, and the number of walls [5–7]. The poor electrical conductivity has imposed great limitations for electronic application of BNNTs. To overcome this obstacle, the most popular method is to modify its electronic properties and to enhance its solubility by chemical functionalization. Many significant efforts have been made for this issue [3, 8–12]. For example, Xie and co-workers have experimentally obtained a soluble BNNT via functionalization with amino-based compounds [8]. Zhi et al. have connected a long alkyl chain to a BNNT via a successful chemical reaction, making the functionalized tubes be soluble in many solvents [11]. Theoretically [13–31], H [13–16], F [17–22], metal [23], NH<sub>3</sub> [24, 25], CCl<sub>2</sub> molecules [26–28], and noncovalent functionalization with organic

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**Fig. 1** Various adsorption sites considered for  $\text{Pd}_3\text{M}$  cluster encapsulation inside the (10, 0) BNNT

molecules [31, 32] have been shown to modify the properties of BNNTs. These studies on the chemical functionalization of BNNTs are important to provide a guidance to design BNNTs-based devices.

In particular, BNNTs filling with molecules/clusters and continuous nanowires made of metals or inorganic compounds may become an interesting and practically important field. This is because the BNNTs are especially suitable for the task of shielding the metallic nanowires inside their cavity to protect the encapsulated content from oxidation during fabrication and application. More importantly, the presence of impurity metal atoms and clusters inside BNNTs during their synthesis can significantly modify the physico-chemical properties and that can differ considerably on the geometry and confinement of the host material. To understand this effect of confinement many experimental and theoretical studies have reported on metal encapsulated BN nanomaterials [33–41]. Zhi and co-workers have summarized filled various materials in BNNTs, including metals and semiconductors [4]. These filled BNNTs can be rather useful. For example, metal-filled BNNTs can be nanoscale “lacquered” wire [33] or natural

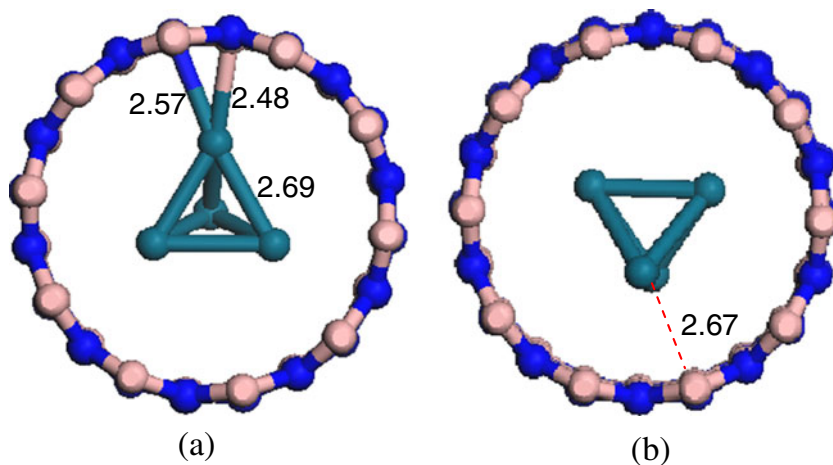
“pipelines” for delivery of tiny metallic clusters under applied thermal, magnetic, or electric fields [34].  $\text{C}_{60}$ -filled BNNTs can work as a special nanoswitch [33]. The interesting electron field-emission of SiC-filled BNNTs can be found [35] and  $\text{MgO}_2$ -filled BNNTs can be used as oxygen generator [36]. Theoretically, the effects of the encapsulation of Fe, Co, Ni, and Cu nanowires on the properties of BNNTs have been reported [37–44]. For example, Xiang et al. have shown that the Ni encapsulated (9, 0) BNNT exhibits interesting semi-half-metallic behavior [37], while Ghosh and co-workers have underscored the possibility of functionalizing these encapsulated BNNTs by interacting them with oxygen molecules [44].

Compared with these above studies, we note that the theoretical reports on the alloys filling in BNNTs are considerable lacking, which is very important not only for deeply understanding the properties of BNNTs, but also for further widening the application fields of BNNTs. Here, using density functional theory (DFT), we have explored the encapsulation of a series of  $\text{Pd}_3\text{M}$  alloys inside a (10, 0) BNNT as an example, where M is Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn, Pt, Pd, or Au, respectively. In particular, the following questions would be mainly addressed: (1) can these  $\text{Pd}_3\text{M}$  stably exist inside the (10, 0) BNNT? (2) If yes, how do the electronic properties of (10, 0) BNNT modify? (3) Whether these  $\text{Pd}_3\text{M}$  clusters can retain their magnetic properties inside BNNT? Apart from these fundamental aspects, one of the primary objectives of this study is to investigate the chemical reactivity of these filled BNNTs by  $\text{Pd}_3\text{M}$  toward oxygen molecule.

### Theoretical methods and models

In the present work, we used DFT methods, implemented in the Dmol<sup>3</sup> package [45, 46], to study the encapsulation of  $\text{Pd}_3\text{M}$  cluster inside the BNNT. All-electrons calculations

**Fig. 2** The obtained stable configurations of  $\text{Pd}_4$  cluster encapsulated inside the (10, 0) BNNT: **a** Config\_1 and **b** Config\_2. The unit of bond length is Å



were employed with the double numerical basis sets plus polarization functional (DNP). The generalized-gradient approximation (GGA) with the Perdew-Burke-Ernzerhof (PBE) [47] and the local density approximation (LDA) with the Perdew-Wang (PWC) functional [48] were applied to estimate a range of the binding energy. In fact, LDA has been shown to be a reliable functional to study the systems involving the van der Waals interactions [49–51] and can give an adsorption energy much closer to the MP2 calculation [52–54]. Additionally, the PBE/DNP method has also been used widely in the study of weak interactions [55–58]. In the present work, GGA with the PBE method is utilized as the exchange-correlation functional throughout the paper. For the 5*d* transition metal atoms Pd, Pt, and Au, the scalar relativistic effect (DSSP) was considered when dealing with their core electrons. The Hirshfeld population analysis [59] was used to obtain both the charges and the net spin populations on each atom. Spin-unrestricted DFT calculations were carried out for Pd<sub>3</sub>M cluster filled in the (10, 0) BNNT with a diameter of 7.94 Å in a periodically repeating tetragonal super cell, whose lattice constants is *a*=*b*=20 Å and *c* (*c* is taken to be twice the one-dimensional lattice parameter of the pure BNNT). The Brillouin zone of the super cell is sampled by 1 × 1 × 3*k* points within the Monkhorst-pack scheme [60]. It should be mentioned that we do not consider the encapsulation of Pd<sub>3</sub>M cluster inside other BNNTs. This is because, for those BNNTs with larger diameters, the encapsulated molecules move almost freely along the nanotube, and eventually leave it. On the contrary, due to being compressed inside BNNTs with smaller diameters, the filled molecules may be excessively compressed, leading to a large structural change for these filled clusters. Thus, they may not stably exist.

### Results and discussion

To determine the stability of the encapsulated Pd<sub>3</sub>M (*M*=Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn, Pd, Pt, and Au) cluster in the BNNT, we first calculated the binding energy, *E<sub>b</sub>*, which is defined as the difference between the total energy of the Pd<sub>3</sub>M clusters encapsulated in the BNNT (*E<sub>Pd3M@BNNT</sub>*), and the sum of the total energies of the (10, 0) BNNT (*E<sub>BNNT</sub>*), and the isolated Pd<sub>3</sub>M cluster (*E<sub>Pd3M</sub>*): *E<sub>b</sub>*=*E<sub>Pd3M@BNNT</sub>*−*E<sub>Pd3M</sub>*−*E<sub>BNNT</sub>*. For each type of the adsorbate, five different configurations were selected to examine the Pd<sub>3</sub>M encapsulation into (10, 0) BNNT, i.e., (1) Pd atop, (2) M atop, (3) Pd–M bridge, (4) Pd–Pd bridge, and (5) Pd–M–Pd face, as shown in Fig. 1, are attached to the inner wall of (10, 0) BNNT, respectively.

Using the Pd<sub>4</sub> case as an example (Fig. 2), for which two stable configurations are obtained: (1) Pd<sub>4</sub> cluster is attached to one side of (10, 0) BNNT, as shown in Fig. 2a (labeled as Config–1). (2) The Pd<sub>4</sub> cluster is locating at the center of the

(10, 0) BNNT (Fig. 2b, Config–2). In the former configuration, the shortest B–Pd and N–Pd bond lengths are 2.48 and 2.57 Å, respectively. For the Config–2, one Pd–Pd bond is parallel to the tube axis, while the other one is perpendicular to the tube axis. The nearest distance between Pd<sub>4</sub> cluster and BNNT is about 2.67 Å. Moreover, the symmetry of the Pd<sub>4</sub> clusters in the above two structures change little after being coated by (10, 0) BNNT. After being inserted, the Pd<sub>4</sub> cluster expands a little, e.g., the average Pd–Pd distance in a free Pd<sub>4</sub> cluster changes from 2.63 to 2.69 Å in Pd<sub>4</sub>@(10, 0) BNNT, and the nanotube walls expand slightly, 0.02 Å for the (10, 0) BNNT. The calculated binding energies for Config–1 and 2 are −1.63 and −1.49 eV, respectively. The negative binding energy value indicates that the encapsulation process of Pd<sub>4</sub> cluster into (10, 0) BNNT is exothermic.

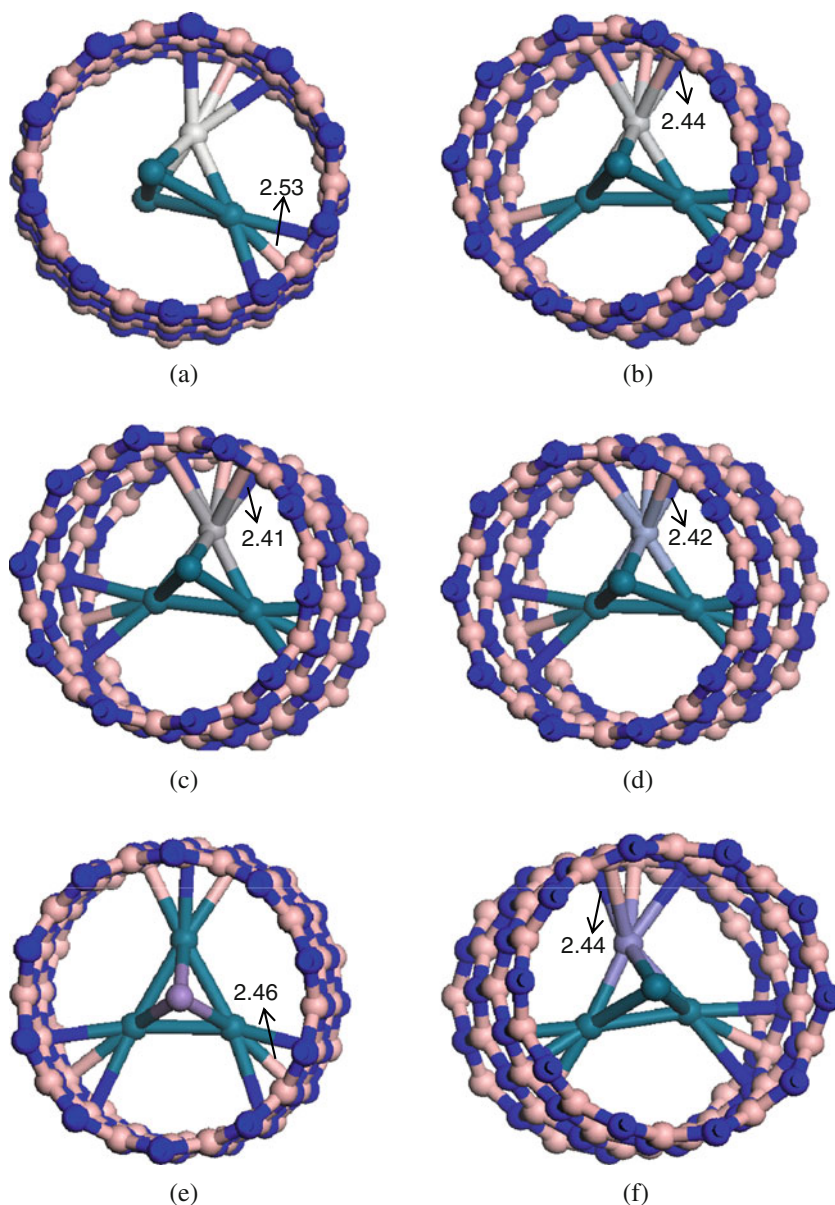
Among the other 12 Pd<sub>3</sub>M clusters considered, the LDA and GGA predict that they can be stably encapsulated inside the (10, 0) BNNT. As expected, the LDA gives a slightly stronger interaction between the Pd<sub>3</sub>M clusters and BNNT. We note that the calculated binding energies range from −0.96 (GGA) to −5.31 eV (GGA), for Pd<sub>3</sub>Sc and Pd<sub>3</sub>V, respectively, as shown in Table 1. The variation of the shortest distance between Pd<sub>3</sub>M and BNNT is consistent with the binding energy, ranging from 2.41 (for Pd<sub>3</sub>V) to 2.53 Å (for Pd<sub>3</sub>Sc). In summary, based on the calculated binding energies, we speculate that those Pd<sub>3</sub>M clusters can be stably filled in the (10, 0) BNNT. Correspondingly, the obtained configurations of Pd<sub>3</sub>M@BNNTs are presented in Fig. 3.

**Table 1** Calculated binding energy (*E<sub>b</sub>*, eV) of a Pd<sub>3</sub>M cluster encapsulated inside the (10, 0) BNNT, the corresponding shortest distance between Pd<sub>3</sub>M cluster and BNNT (*d*, Å), the charge transfers from the Pd<sub>3</sub>M cluster to BNNT (*Q*, *e*), and the net magnetic moments of the M atom/total system

M	<i>E<sub>b</sub></i> <sup>a</sup>	<i>d</i>	<i>Q</i>	<i>μ<sub>M</sub></i> / <i>μ<sub>total</sub></i>
Sc	−0.96 (−1.24)	2.53	0.33	0
Ti	−1.81 (−2.32)	2.44	0.39	1.44/1.84
V	−5.31 (−5.92)	2.41	0.36	2.59/2.99
Cr	−4.67 (−5.16)	2.42	0.42	3.77/4.09
Mn	−2.97 (−3.31)	2.46	0.63	5.01/5.08
Fe	−2.89 (−3.14)	2.44	0.37	3.81/3.91
Co	−1.69 (−2.20)	2.46	0.59	2.90/2.94
Ni	−1.78 (−2.04)	2.52	0.37	1.86/1.91
Cu	−1.70 (−1.97)	2.46	0.53	0.43/0.44
Zn	−1.43 (−1.75)	2.46	0.61	0
Pd	−1.63 (−1.82)	2.48	0.56	1.49/1.58
Pt	−1.82 (−2.01)	2.45	0.48	1.32/1.40
Au	−1.77 (−1.95)	2.46	0.52	0.04/0.05

<sup>a</sup> Values in parentheses are derived from the LDA/PWC method

**Fig. 3** The obtained most stable configurations of Pd<sub>3</sub>M cluster encapsulated inside the (10, 0) BNNT, where M is **a** Sc, **b** Ti, **c** V, **d** Cr, **e** Mn, **f** Fe, **g** Co, **h** Ni, **i** Cu, **j** Zn, **k** Pt, and **l** Au. The unit of bond length is Å

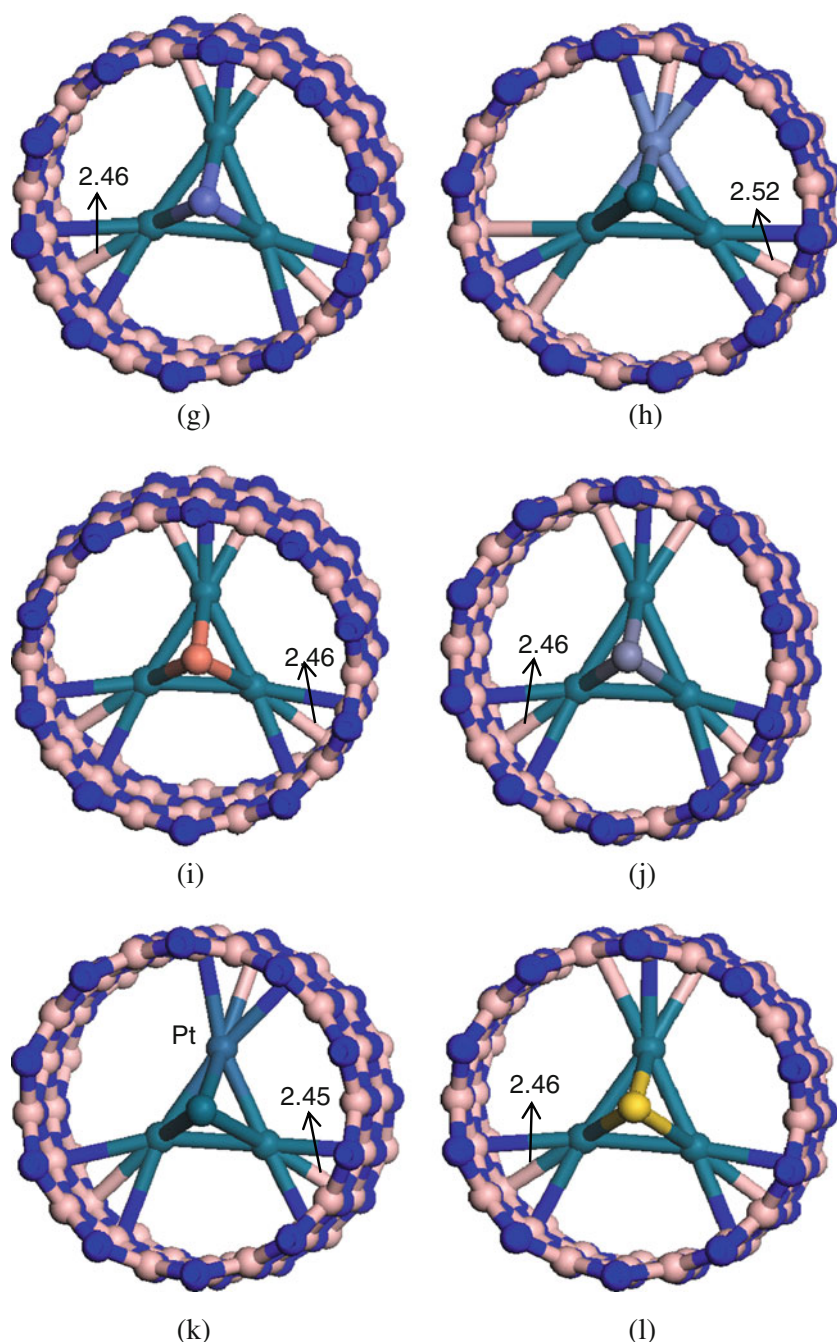


In addition to the binding energies and local-minimum configurations, we have examined the effects of Pd<sub>3</sub>M clusters encapsulation on the electronic structures of (10, 0) BNNT. Based on the spin unrestricted calculations, the band structures of the pristine (10, 0) BNNT as well as those of the most stable 13 Pd<sub>3</sub>M@BNNT systems are calculated. Both the plus and minus spins (labeled with “+” and “-” in Fig. 4) were considered. For the pristine BNNT, the band structures for the plus and minus spins are the same, exhibiting a direct band gap of 4.13 eV (Fig. 5a). Moreover, it can be seen from Fig. 5 that in all 13 systems the encapsulation of the Pd<sub>3</sub>M clusters induces certain impurity states within the band gap of the pristine BNNT, thereby leading to the band gap reduction. In most cases, the band gap is no greater than 1.0 eV. The results of projected density of states (PDOS) indicate that these impurity states mainly due to

the *d* electrons of the Pd<sub>3</sub>M clusters. To some extent, the *p* electrons also contribute to the density of states near the Fermi level, but the contribution is much less than the *d* electrons. The BNNT contributes little to the density of states near the Fermi level.

The electron-charge analysis using the Hirshfeld method is summarized in Table 1. With the encapsulation of Pd<sub>3</sub>M clusters inside BNNTs, certain amount charge is transferred from Pd<sub>3</sub>M clusters to BNNT, ranging from 0.33 *e* (Pd<sub>3</sub>Sc) to 0.63 *e* (Pd<sub>3</sub>Mn). As a result, a net magnetic moment may emerge. Because the ground state of the pristine BNNT is nonmagnetic, the net spin mainly originates from the magnetism of Pd<sub>3</sub>M clusters as shown in Table 1. Moreover, Pd<sub>3</sub>M encapsulation also leads to an induction of local magnetic moment on the N atoms. This is expected, because the 3*d* orbitals of Pd<sub>3</sub>M clusters gets partially hybridized with N 2*p*

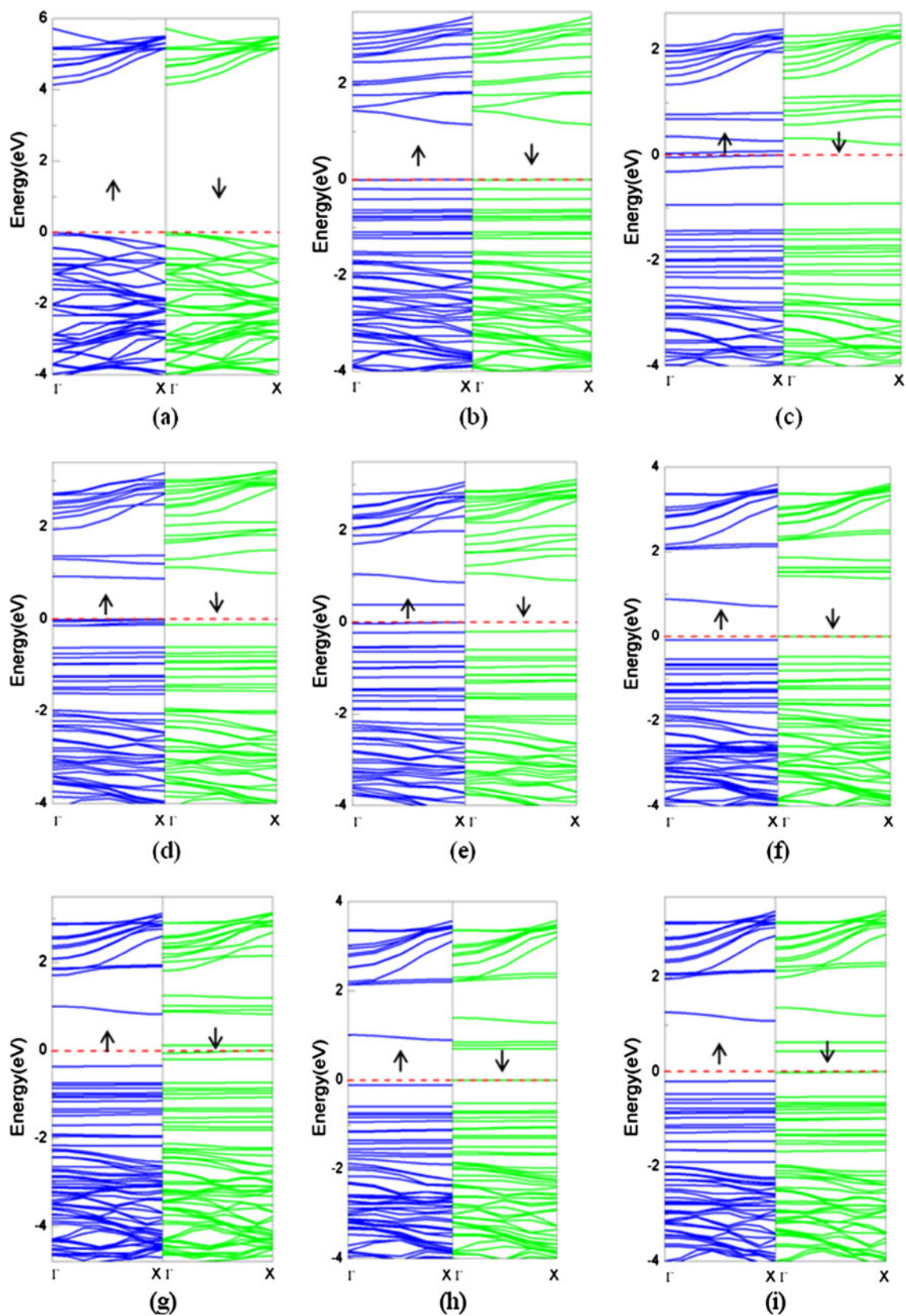
Fig. 3 (continued)



states. Among all  $\text{Pd}_3\text{M}$  clusters that have been studied here,  $\text{Pd}_3\text{Mn}$  shows the largest spin polarization (total magnetic moment of  $5.08 \mu_B$ ), whereas that of  $\text{Pd}_3\text{Au}$  is smallest (total magnetic moment of  $0.05 \mu_B$ ), and the value decreases to 0 for  $\text{Pd}_3\text{Zn}$ . The small magnetic moment of filled BNNT by  $\text{Pd}_3\text{Zn}$  and  $\text{Pd}_3\text{Au}$  seems to correlate with the valence shell of Zn and Au: Zn has a filled valence shell (there is no unpaired electron) and Au has only one unpaired electron in its valence shell. Thus, when the two clusters are filled inside BNNT, their valence electrons involve in interaction with BNNT, decreasing their magnetic moment, compared with their freestanding ones. In terms of the unique

properties, the filled BNNT by some  $\text{Pd}_3\text{M}$  clusters (such as  $\text{Pd}_3\text{Mn}$ ) may have potential applications in spintronic and quantum information.

Due to the appearance of mid-gap states around Fermi level of  $\text{Pd}_3\text{M@BNNT}$  and charge transfer between BNNT and  $\text{Pd}_3\text{M}$  clusters, the chemical reactivity of BNNT may be enhanced to different degrees. To verify this fact, we have studied the adsorption of an  $\text{O}_2$  molecule on  $\text{Pd}_3\text{M@BNNT}$  composites and compared it with that on pristine BNNT. The calculation is performed by placing a single  $\text{O}_2$  molecule on  $\text{Pd}_3\text{M@BNNT}$  in vertical or horizontal orientations at different sites, including atop of B or N, edge of the B–N



**Fig. 4** The band structures of a pristine (10, 0) BNNT, and Pd<sub>3</sub>M@BNNT, where M is b Sc, c Ti, d V, e Cr, f Mn, g Fe, h Co, i Ni, j Cu, k Zn, l Pd, m Pt, and n Au. The plus and minus spin electronic structures are distinguished with “+” and “-”. The Fermi level is plotted with the red dotted line

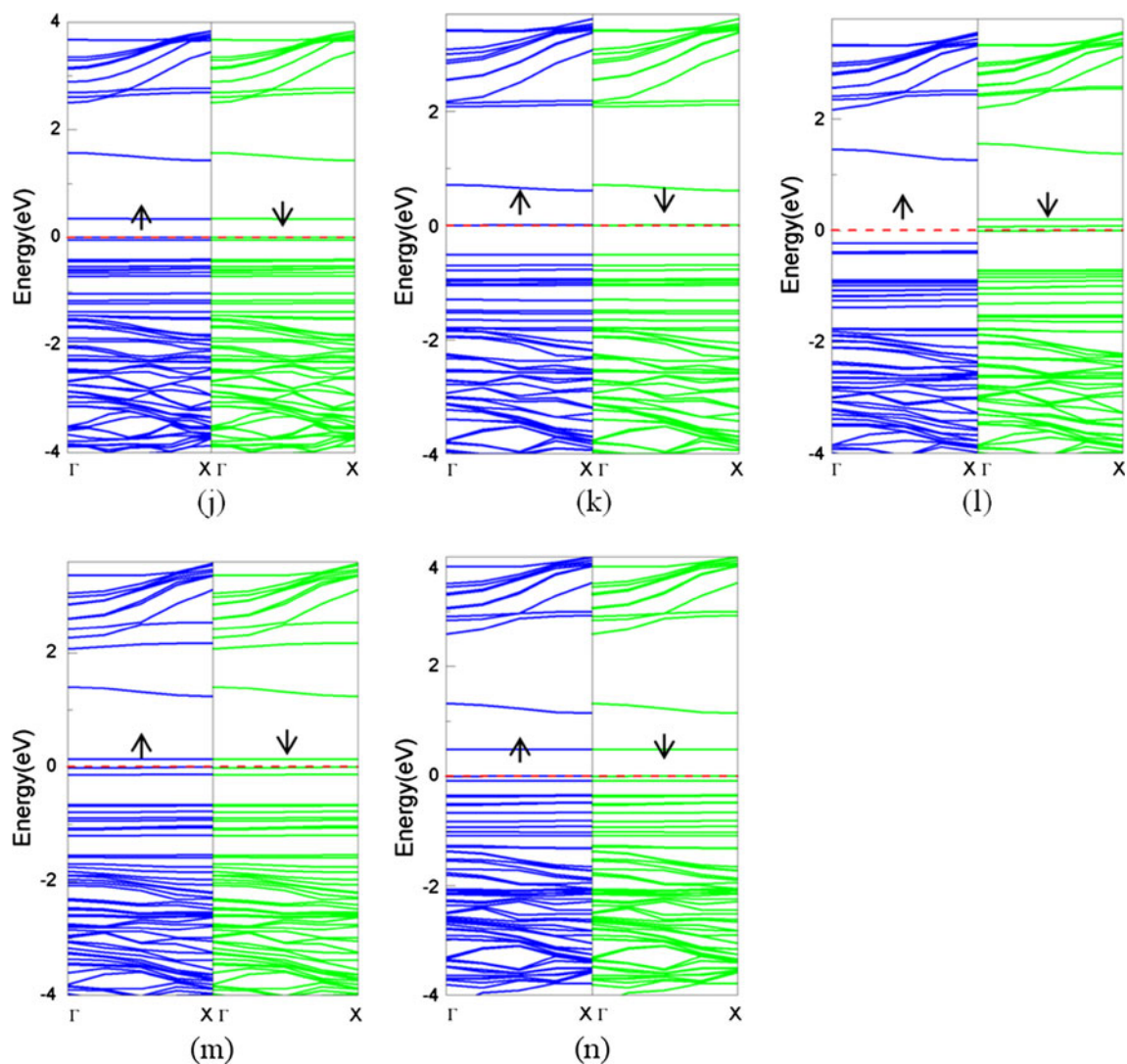
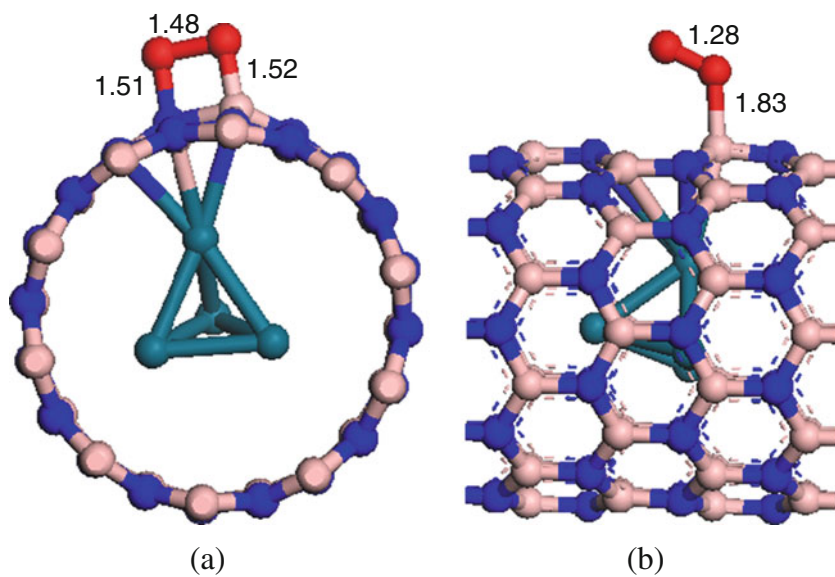


Fig. 4 (continued)

Fig. 5 The optimized stable configurations of a single O<sub>2</sub> molecule adsorbed on Pd<sub>4</sub>@BNNT in (a) horizontal and (b) vertical orientation. The unit of bond length is Å



bond, and hexagon of  $B_3N_3$ . The calculated adsorption energies, structural parameters, and charge transfer of  $O_2$  on  $Pd_3M@BNNTs$  are shown in Table 2. For the  $O_2$  molecule on pristine BNNT, we find that this kind of adsorption is considerably weaker with  $E_{ads}$  of  $-0.14$  eV. The distance between BNNT and the  $O_2$  molecule is  $3.12$  Å.

In contrast, for  $Pd_4@BNNT$ , the  $O_2$  molecule can be stably adsorbed on BNNT. The lowest energy configuration is that  $O_2$  molecule prefers to adsorb in horizontal orientation, as shown in Fig. 5a. For this configuration, the O-O bond is elongated up to  $1.48$  Å (peroxy type [61]), which is larger than in its free form ( $1.22$  Å). Meanwhile, the  $O_2$  molecule binds with the axial B-N bond with B-O and N-O bond lengths of  $1.52$  and  $1.51$  Å, respectively. The energy released during this process is estimated to be  $2.36$  eV and about  $0.22$  electrons are transferred from  $Pd_4@BNNT$  to  $O_2$  molecule. As shown in Fig. 5b, a meta-stable configuration is also obtained for  $O_2$  adsorption on  $Pd_4@BNNT$ , where the  $O_2$  molecule adsorbs on one of B atom of BNNT in elbow shaped orientation. The calculated adsorption energy for this configuration is  $-0.84$  eV. The O-O and O-B bond lengths are  $1.28$  and  $1.83$  Å, respectively.

In Table 2, we present the results of  $O_2$  interactions with other  $Pd_3M@BNNT$  materials. The results suggest that the binding strength of  $O_2$  molecule with these encapsulated BNNTs is much stronger than that with pristine BNNT in various ways. Moreover, the adsorption energy of  $O_2$

molecule on  $Pd_3Sc@BNNT$  is the largest ( $E_{ads} = -2.09$  eV), suggesting a strong chemisorption. Similarly, the encapsulation of  $Pd_3Ti$ ,  $Pd_3V$ ,  $Pd_3Cr$ , and  $Pd_3Fe$  also renders BNNT high chemical reactivity toward  $O_2$  with large adsorption energies of  $-1.62$ ,  $-1.18$ ,  $-0.83$ ,  $-0.68$  eV, respectively. Especially, the O-O bonds of adsorbed  $O_2$  molecules on the above BNNT-based composites are elongated to  $\sim 1.33$  Å. This fact indicates that the filled BNNTs with  $Pd_3Sc$ ,  $Pd_3Ti$ ,  $Pd_3V$ ,  $Pd_3Cr$ , and  $Pd_3Fe$  clusters have potential for developing oxidative catalysis. For other  $Pd_3M@BNNTs$  ( $M=Mn, Co, Ni, Cu, Zn, Pt, \text{ and } Au$ ), however, we find that the chemical reactivity of BNNT toward  $O_2$  molecule ( $E_{ads} = -0.27, -0.27, -0.38, -0.22, -0.27, -0.22, \text{ and } -0.23$  eV, respectively) is slightly stronger than the pristine one ( $E_{ads} = -0.14$  eV). This can be testified by the large distance between BNNT and adsorbates as shown in Table 2. Due to the weak interaction, it is expected that the seven  $Pd_3M@BNNTs$  ( $M=Mn, Co, Ni, Cu, Zn, Pt, \text{ and } Au$ ) are unsuitable for oxidative catalysis. On the basis of the above results, it is found that the binding strength of  $O_2$  with  $Pd_3M@BNNTs$  seems to be unrelated with the magnetic momentums of these nanocomposites.

## Conclusions

In summary, we have carried out density functional theory study to calculate the equilibrium structures and energetics of different small  $Pd_3M$  clusters ( $M=Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn, Pd, Pt, \text{ and } Au$ ) encapsulated in (10, 0) BNNT. In particular, the effects of the  $Pd_3M$ -encapsulation on the electronic and magnetic properties are mainly addressed. The results indicate that these considered  $Pd_3M$  clusters can be stably filled in (10, 0) BNNT. The band structures of  $Pd_3M@BNNTs$  show that some new impurity bands are introduced to the large band gap of the pristine BNNT, rendering it to exhibit semiconducting nature, whose band gaps are less than  $1.0$  eV. Moreover, the encapsulation of the  $Pd_3M$  clusters can give rise to a variety of net magnetic moments, ranging from  $5.08 \mu_B$  to  $0$ . Finally, by studying the interaction of BNNT and  $O_2$  molecules, we find that the encapsulated BNNTs by  $Pd_3M$  cluster ( $M=Sc, Ti, V, Cr, Fe, \text{ and } Pd$ ) exhibit a much higher chemical reactivity than the pristine one. The elongation of the O-O bonds of adsorbed  $O_2$  molecules indicates that several  $Pd_3M@BNNTs$  can also be used for further industrial catalytic applications. Though the current DMol3 based DFT-GGA calculations are not very accurate in terms of the absolute magnitude of band gap, we expect that our qualitative results on the effects of  $Pd_3M$  encapsulation on the electronic properties of the BNNT remain robust with respect to more accurate calculations.

**Table 2** The calculated adsorption energy ( $E_{ads}$ , eV)<sup>a</sup> of the  $O_2$  molecule on  $Pd_3M@BNNT$ , charge transfer from  $Pd_3M@BNNT$  to  $O_2$  molecule ( $Q$ , e), optimized O-O ( $d_{O-O}$ ), and B-O distances ( $d_{B-O}$ , Å). The O-O bond length in the free  $O_2$  molecule is calculated to be  $1.22$  Å

M	$E_{ads}$	$Q$	$d_{B-O}$	$d_{O-O}$
pure BNNT	-0.14	-0.01	3.12	1.23
Sc	-2.09	-0.27	1.50	1.35
Ti	-1.62	-0.27	1.51	1.35
V	-1.18	-0.26	1.50	1.35
Cr	-0.83	-0.23	1.56	1.33
Mn	-0.27	-0.13	3.12	1.24
Fe	-0.68	-0.22	1.56	1.33
Co	-0.27	-0.13	2.87	1.24
Ni	-0.38	-0.16	2.83	1.25
Cu	-0.22	-0.11	3.08	1.24
Pd	-2.36	-0.22	1.52	1.48
Zn	-0.27	-0.13	3.03	1.24
Pt	-0.22	-0.12	2.92	1.24
Au	-0.23	-0.11	3.13	1.24

<sup>a</sup> The adsorption energy is defined as:  $E_{ads} = E_{total}(O_2 + Pd_3M@BNNT) - E_{total}(Pd_3M@BNNT) - E_{total}(O_2)$ , where  $E_{total}$  is the total energy of the system in the bracket



**Acknowledgments** This work is supported by the National Nature Science Foundation of China (No. 21203048) and University Key Teacher Foundation of Heilongjiang Provincial Education Department (No. 1252G030).

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